

pro $\frac{1}{\sqrt{2}} (|0\rangle + |2\rangle)$

$E = \frac{1}{2} \left(\frac{1}{2} \hbar \omega + \frac{5}{2} \hbar \omega \right) = \frac{3}{2} \hbar \omega \rightarrow \frac{\partial E}{\partial \omega} = \frac{3}{2} \hbar$

$\langle \psi | m \omega x^2 | \psi \rangle = \frac{1}{2} (\langle 0 | + \langle 2 |) \left[\frac{\hbar}{2} (q^2 + q^{+2} + 2q^+q + 1) \right] (|0\rangle + |2\rangle)$
 $= \frac{\hbar}{4} \left[\langle 0 | 2q^+q + 1 | 0 \rangle + \langle 2 | 2q^+q + 1 | 2 \rangle + \langle 0 | q^2 | 2 \rangle + \langle 2 | q^{+2} | 0 \rangle \right]$
 $\hookrightarrow = \langle 0 | q^2 | 2 \rangle = \langle 0 | \sqrt{2} | 0 \rangle = \sqrt{2} = \langle 2 | q^{+2} | 0 \rangle = \sqrt{2}$
 $= \frac{\hbar}{4} [1 + 3 \cdot 5 + 2\sqrt{2}] = \frac{3}{2} \hbar \omega + \frac{\hbar}{\sqrt{2}}$ korekci s patne b

Mejme LHO ^{casnici stavu} $|1\rangle = \frac{1}{\sqrt{3}} (|0\rangle + |1\rangle + |3\rangle)$

overle $\frac{d\langle A \rangle}{dt} = \frac{i}{\hbar} \langle [H, A] \rangle + \left\langle \frac{\partial A}{\partial t} \right\rangle$

pro $\hat{A} = \hat{p}$

$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$

$[H, p] = \left[\frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2, p \right]$

$= \left[\frac{p^2}{2m}, p \right] + \frac{1}{2} m \omega^2 [x^2, p]$

$= 0 + \frac{1}{2} m \omega^2 [x^2 p - p x^2],$ z q'ar $[x, p] = i\hbar$

$x^2 p - p x^2 = x x p - p x x = x x p - x p x + x p x - p x x$
 $\downarrow \qquad \qquad \qquad \downarrow$
 $x [x, p] \qquad [x, p] x$

$= x i\hbar + x i\hbar = 2i\hbar x$

$[H, p] = \frac{1}{2} m \omega^2 2i\hbar x = i m \omega^2 \hbar x$

$= i m \omega^2 \hbar \frac{d}{dt} (q + q^+)$

$p = \frac{\sqrt{\hbar m \omega}}{\sqrt{2} i} (q - q^+)$

$\langle p \rangle = \frac{1}{3} (\langle 0 | + \langle 1 | + \langle 3 |) \frac{\sqrt{\hbar m \omega}}{\sqrt{2} i} (q - q^+) (|0\rangle + |1\rangle + |3\rangle)$
 $= \frac{1}{3} \frac{\sqrt{\hbar m \omega}}{\sqrt{2} i} \left[\langle 0 | q | 1 \rangle e^{-i\omega t} - \langle 1 | q^+ | 0 \rangle e^{i\omega t} \right]$

$\frac{\partial \langle A \rangle}{\partial t} = -\frac{\sqrt{2}}{3} \sqrt{\hbar m \omega} \omega \cos(\omega t)$

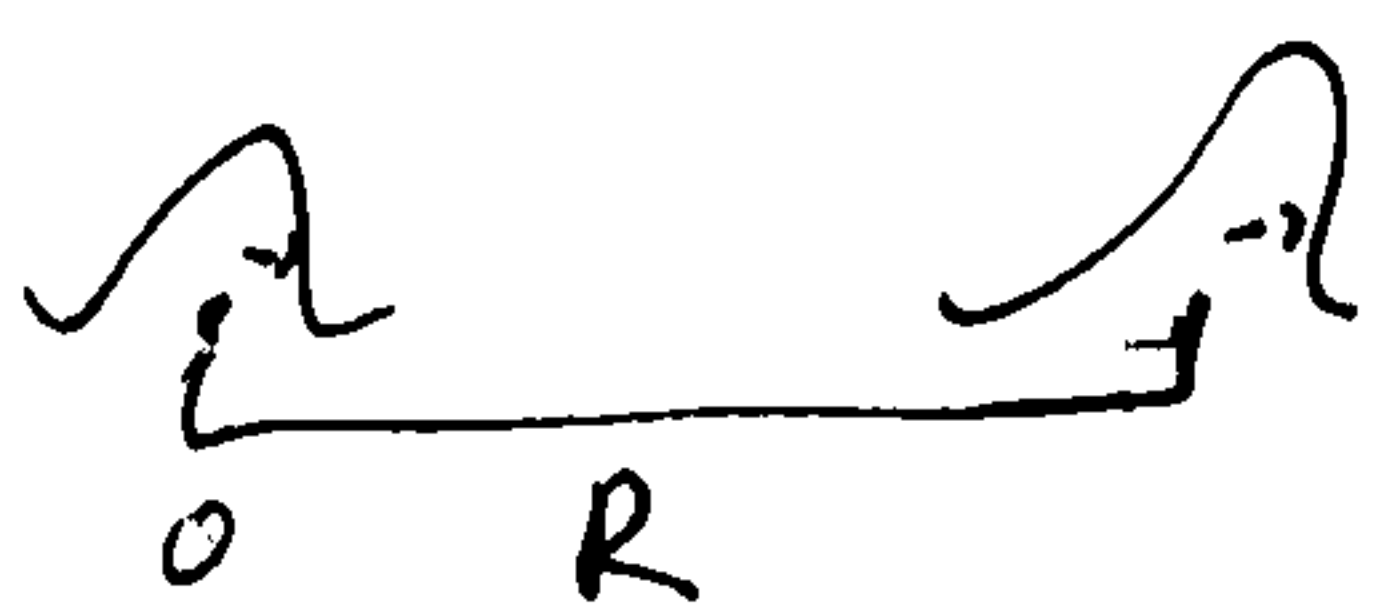
$$\begin{aligned} \frac{d\langle A \rangle}{dt} &= \frac{i}{\hbar} \langle [H, A] \rangle + \left\langle \frac{\partial A}{\partial t} \right\rangle \\ &\rightarrow \phi \\ &= \frac{i}{\hbar} \langle i m \omega^2 \frac{x}{\sqrt{2}} (a + a^\dagger) \rangle \\ &= \frac{i}{\hbar} - m \omega^2 \sqrt{\frac{\hbar}{m\omega}} \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} \left[\langle 0 | e^{\frac{i\omega t}{2}} + \langle 1 | e^{\frac{3i\omega t}{2}} + \langle 3 | e^{\frac{7i\omega t}{2}} \right] \\ &\quad (a + a^\dagger) \left[|0\rangle e^{-\frac{i\omega t}{2}} + |1\rangle e^{-\frac{3i\omega t}{2}} + |3\rangle e^{-\frac{7i\omega t}{2}} \right] \\ &= -\frac{1}{3} \sqrt{\frac{\hbar}{m\omega}} \frac{\omega}{\sqrt{2}} \left[\langle 0 | a | 1 \rangle e^{-i\omega t} + \langle 1 | a^\dagger | 0 \rangle e^{i\omega t} \right] \\ &\quad \underbrace{\hspace{10em}}_{2 \cos(\omega t)} \\ &= -\frac{\sqrt{2}}{3} \sqrt{\frac{\hbar}{m\omega}} \omega \cos(\omega t) \end{aligned}$$

$$\begin{aligned} \frac{dE_\lambda}{d\lambda} &= \frac{d \langle \psi_\lambda | H_\lambda | \psi_\lambda \rangle}{d\lambda} \\ &= \left\langle \frac{d\psi_\lambda}{d\lambda} | H_\lambda | \psi_\lambda \right\rangle + \underbrace{\langle \psi_\lambda | \frac{dH_\lambda}{d\lambda} | \psi_\lambda \rangle}_{\text{v.l. step}} + \left\langle \psi_\lambda | H_\lambda | \frac{d\psi_\lambda}{d\lambda} \right\rangle \\ &\quad \left\langle \frac{d\psi_\lambda}{d\lambda} | E_\lambda | \psi_\lambda \right\rangle \quad \leftarrow \quad E_\lambda \left\langle \psi_\lambda | \frac{d\psi_\lambda}{d\lambda} \right\rangle \\ &\quad \leftarrow \quad E_\lambda \frac{d \langle \psi_\lambda | \psi_\lambda \rangle}{d\lambda} = 0 \\ &= \left\langle \psi_\lambda | \frac{dH_\lambda}{d\lambda} | \psi_\lambda \right\rangle \end{aligned}$$

$$\begin{aligned} |\phi\rangle &= \frac{1}{\sqrt{\alpha} \sqrt{\pi}} e^{-\frac{x^2}{2\alpha^2}} \quad \alpha = \sqrt{\frac{\hbar}{m\omega}} \quad \frac{d\alpha}{d\omega} = \\ \frac{d|\phi\rangle}{d\omega} &= \frac{d}{d\omega} \left[\frac{1}{\pi^{1/4}} \left(\frac{m\omega}{\hbar}\right)^{1/4} e^{-\frac{x^2 m \omega}{2\hbar}} \right] = \left(\frac{m}{\pi \hbar}\right)^{1/4} \left[\frac{1}{4} \omega^{-3/4} - \frac{x^2 m}{2\hbar} \omega^{1/4} \right] e^{-\frac{x^2 m \omega}{2\hbar}} \end{aligned}$$

Send pro $\alpha = \frac{\hbar}{m\omega} = 1$

• Májme 2 LHO ve vzdálenosti R



• částice realizované

$$|\psi_1^0\rangle = |\phi_1\rangle \quad |\psi_2^0\rangle = |\phi_2\rangle$$

Vlnová funkce celku

$$|\psi\rangle = |\phi_1\rangle \otimes |\phi_2\rangle = |\phi\phi\rangle$$

$$= \frac{1}{\sqrt{\alpha\sqrt{\pi}}} e^{-\frac{x_1^2}{2\alpha^2}} \cdot \frac{1}{\sqrt{\alpha\sqrt{\pi}}} e^{-\frac{(x_2-R)^2}{2\alpha^2}}$$

• Energie jednotlivé & celkové

$$E_1 = \hbar\omega(n_1 + 1/2) \quad E_2 = \hbar\omega(n_2 + 1/2)$$

$$E_{\text{tot}} = \hbar\omega(n_1 + n_2 + 1)$$

• ~~total~~ a, a[†] operatory:

$$1: a, a^\dagger \quad 2: b, b^\dagger \quad , \quad x_1 = \frac{\alpha}{\sqrt{2}}(a + a^\dagger) \quad x_2 = \frac{\alpha}{\sqrt{2}}(b + b^\dagger)$$

$$\text{tot} \quad a \otimes 1, a^\dagger \otimes 1 \quad 1 \otimes b, 1 \otimes b^\dagger$$

$$\text{působí} \quad a \otimes 1 |\phi\phi\rangle = 0, \quad a^\dagger \otimes 1 |\phi\phi\rangle = \sqrt{1} |1\phi\rangle$$

Uvažujme, že částice jsou nabité a "připojené" pružinami k ke jádru v ϕ , resp. R (Drudeho model)

$$3D: \quad \vec{r}_1, \vec{r}_2 \quad V = \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{|R|} + \frac{1}{|R - \vec{r}_1 + \vec{r}_2|} - \frac{1}{|R - \vec{r}_1|} - \frac{1}{|R - \vec{r}_2|} \right)$$

di approx $|\vec{r}_1|, |\vec{r}_2| \ll R$ ($|R - \vec{r}_1| = \sqrt{(R_x - r_{1x})^2 + (R_y - r_{1y})^2 + (R_z - r_{1z})^2}$)

$$V = \frac{e^2}{4\pi\epsilon_0} \left(\frac{\vec{r}_1 \cdot \vec{r}_2}{R^3} - \frac{3(\vec{R} \cdot \vec{r}_1)(\vec{R} \cdot \vec{r}_2)}{R^5} \right)$$

Taylor $(1 - \frac{R \cdot r}{R^2})^{-1/2}$

• Uvažujme $\vec{R} = (R, \phi, \theta)$

$$\rightarrow V = \frac{e^2}{4\pi\epsilon_0} \left(\frac{\vec{r}_1 \cdot \vec{r}_2}{R^3} - 3 \frac{R_x x_1 x_2 + y_1 y_2 + z_1 z_2}{R^5} - \frac{3 x_1 x_2}{R^3} \right)$$

$$= \frac{e^2}{4\pi\epsilon_0} \left(-\frac{2x_1 x_2}{R^3} + \frac{y_1 y_2 + z_1 z_2}{R^3} \right)$$

~~Va~~ $V = \frac{e^2}{4\pi\epsilon_0} \left(-\frac{2x_1 x_2}{R^3} \right)$ Uvažujme jón x

$$V = \frac{e^2}{4\pi\epsilon_0} \left(-\frac{2x_1 x_2}{R^3} \right)$$

zapíšeme pomocí $a, a^\dagger, \downarrow, \downarrow^\dagger$

$$= -\frac{2e^2}{4\pi\epsilon_0 R^3} (a + a^\dagger)(\downarrow + \downarrow^\dagger) \frac{d^2}{2}$$

$$= -\frac{e^2 d^2}{4\pi\epsilon_0 R^3} (a + a^\dagger)(a \otimes \downarrow + a \otimes \downarrow^\dagger + a^\dagger \otimes \downarrow + a^\dagger \otimes \downarrow^\dagger)$$

Matrixové elementy se základem stavem

$$a \otimes \downarrow |00\rangle = 0$$

$$a \otimes \downarrow^\dagger |00\rangle = 0$$

$$a^\dagger \otimes \downarrow |00\rangle = 0$$

$$a^\dagger \otimes \downarrow^\dagger |00\rangle = |11\rangle = |11\rangle$$

$$\rightarrow \begin{array}{c|cc} & |00\rangle & |11\rangle \\ \hline \langle 00| & \frac{1}{2}\omega & -K \\ \langle 11| & -K & 3\frac{1}{2}\omega \end{array}$$

$$K = \frac{e^2 d^2}{4\pi\epsilon_0 R^3}$$

$$H = \begin{array}{c|cc} & |00\rangle & |11\rangle \\ \hline \langle 00| & \frac{1}{2}\omega & -K \\ \langle 11| & -K & 3\frac{1}{2}\omega \end{array}$$

$$\begin{vmatrix} \frac{1}{2}\omega - \lambda & -K \\ -K & 3\frac{1}{2}\omega - \lambda \end{vmatrix} = (\frac{1}{2}\omega - \lambda)(3\frac{1}{2}\omega - \lambda) - K^2 = \lambda^2 - 4\frac{1}{2}\omega\lambda + 3\frac{1}{2}^2\omega^2 - K^2$$

$$D = B^2 - 4AC = 16\frac{1}{2}^2\omega^2 - 4 \cdot (3\frac{1}{2}^2\omega^2 - K^2) = 4\frac{1}{2}^2\omega^2 + 4K^2$$

$$\lambda_{1,2} = \frac{-B \pm \sqrt{D}}{2A} = \frac{4\frac{1}{2}\omega \pm \sqrt{4\frac{1}{2}^2\omega^2 + 4K^2}}{2} = 2\frac{1}{2}\omega \pm \sqrt{\frac{1}{2}^2\omega^2 + K^2}$$

$$= 2\frac{1}{2}\omega \pm \frac{1}{2}\omega \sqrt{1 + \frac{K^2}{\frac{1}{2}^2\omega^2}} = 2\frac{1}{2}\omega \pm \frac{1}{2}\omega \left(1 + \frac{1}{2} \frac{K^2}{\frac{1}{2}^2\omega^2} \right)$$

$$\pm \lambda_+ = 3\frac{1}{2}\omega + \frac{1}{2} \frac{K^2}{\frac{1}{2}\omega}$$

$$\lambda_- = \frac{1}{2}\omega - \frac{1}{2} \frac{K^2}{\frac{1}{2}\omega} = \frac{1}{2}\omega - \frac{1}{2} \frac{e^4 d^4}{(4\pi\epsilon_0)^2 R^6} \frac{1}{\frac{1}{2}\omega}$$

$$d^4 = \left(\frac{\hbar}{m\omega} \right)^4 = \frac{\hbar^4}{m^2\omega^2}$$

$$= \frac{1}{2}\omega - \frac{1}{2} \frac{e^4}{(4\pi\epsilon_0)^2} \frac{\hbar^4}{m^2\omega^3} \frac{1}{R^6}$$

$-\frac{1}{R^6} \sim$ vdW inter.

• LHO - opisani, v dlu, x^2 časovni nivoj, a, a' na x-ostev, a, a' na p-ostev, OC, HF theorem, x^3 v normalni ordeni, x^2 časovni nivoj, a, a' na x-ostev, a, a' na p-ostev, OC, HF theorem, x^3 v normalni ordeni, x^2 časovni nivoj, a, a' na x-ostev, a, a' na p-ostev, OC, HF theorem, x^3 v normalni ordeni

• Mejnme $\psi = \frac{2}{\sqrt{3d\sqrt{\pi}}} \frac{x^2}{d^2} e^{-\frac{x^2}{2d^2}}$, jaky' je časovni nivoj? pokud se nahajati v potencialu $\frac{1}{2} m\omega^2 x^2$

• Normalizacno:

$$\int_{-\infty}^{\infty} \frac{4}{3d\sqrt{\pi}} \frac{x^4}{d^4} e^{-\frac{x^2}{d^2}} dx = \frac{4}{3d\sqrt{\pi}} \frac{1}{d^4} \cdot \frac{d^2}{2} \frac{d^2}{2} \sqrt{\pi} d = 1 \quad \text{OK}$$

~~normalizirani~~ $\langle 0 | \psi \rangle = \frac{1}{\sqrt{d\sqrt{\pi}}} \frac{2}{\sqrt{3d\sqrt{\pi}}} \int_{-\infty}^{\infty} \frac{x^2}{d^2} e^{-\frac{x^2}{d^2}} dx$

$$|0\rangle = \frac{1}{\sqrt{d\sqrt{\pi}}} e^{-\frac{x^2}{2d^2}} = \frac{2}{2\sqrt{\pi} \sqrt{3}} \frac{1}{\sqrt{d}} \sqrt{\frac{d^2}{2}} = \frac{1}{\sqrt{3}} = c_0$$

$c_2 = \sqrt{1 - c_0^2} = \sqrt{\frac{2}{3}}$... v principu $e^{i\delta}$ poravnava in s $|0\rangle$ zjeshine, če $\delta = \alpha$

$$\rightarrow |\psi\rangle = \frac{1}{\sqrt{3}} |0\rangle + \sqrt{\frac{2}{3}} |2\rangle$$

$$|\psi\rangle = \frac{1}{\sqrt{3}} |0\rangle e^{-\frac{i\omega t}{2}} + \sqrt{\frac{2}{3}} |2\rangle e^{-\frac{i5\omega t}{2}}$$

$$\rho = \psi^* \psi = \left[\frac{1}{\sqrt{3}} \frac{1}{\sqrt{d\sqrt{\pi}}} e^{-\frac{x^2}{2d^2}} e^{\frac{i\omega t}{2}} + \sqrt{\frac{2}{3}} \frac{1}{\sqrt{2d\sqrt{\pi}}} \left[\frac{2x^2}{d^2} - 1 \right] e^{-\frac{x^2}{2d^2}} e^{\frac{i5\omega t}{2}} \right]$$

$$= \frac{1}{d\sqrt{\pi}} \frac{1}{3} e^{-\frac{x^2}{d^2}} + \frac{2}{3} \frac{1}{2d\sqrt{\pi}} e^{-\frac{x^2}{d^2}} \left[\frac{2x^2}{d^2} - 1 \right]^2 + \frac{\sqrt{2}}{3} \frac{1}{d\sqrt{\pi}} \frac{1}{\sqrt{2}} e^{-\frac{x^2}{d^2}} \left[e^{-2i\omega t} + e^{2i\omega t} \right] \left[\frac{2x^2}{d^2} - 1 \right]$$

$2 \cos(2\omega t) \sim \cos(\Delta E)$

